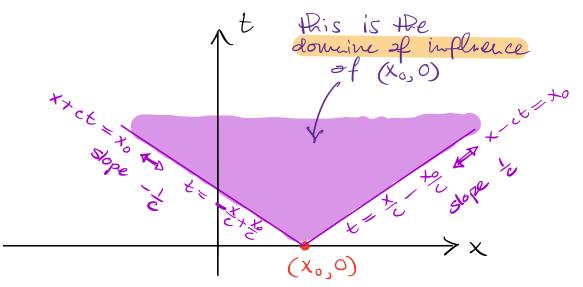
2.2 Consolity and Energy As we have seen, the vave equation in 1D is the sum of two distarbaces: our moving to the right and the other moving to the beft.

> $2 \neq = 0$ \longrightarrow signeds more at exactly C. $2 \neq \neq 0$ \implies signeds more at $\leq C$.

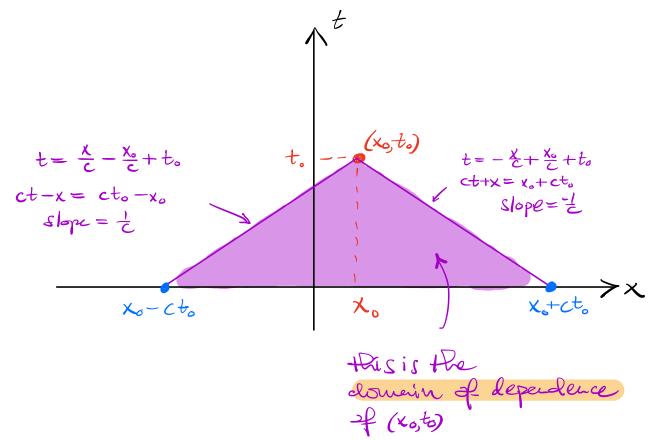
(we havit proven this, but it can be seen by boking at d'Alembert's Jornula, and, Rinking...)

We useably sketch time-dependent problems with t (time) being the verticed axis, and the Porizontal axis representing the spatial coordinate(s) (for now were in 1D, so there's just x; but is principle there can be x, y, z...)

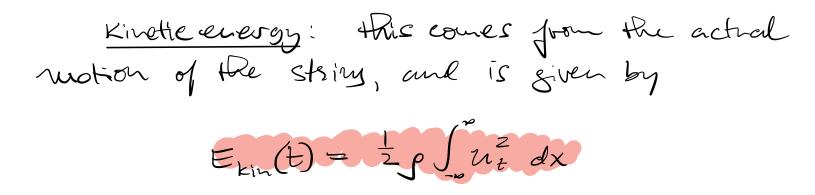


That is:

Take some time
$$s < t_0$$
. Then what are
all the points $y \in \mathbb{R}$ sit. a signal
leaving y at time s will reach x_0
at time to?

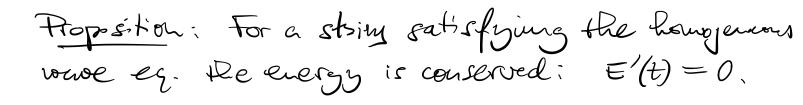


<u>Evergy</u>: There is a natural way to associate a non-vegative grantity we call "energy" to the wave equation on R. We denote the total energy E. It is made up of two parts:



Potential every: this is the energy stored in the string, for example when it is stretched. $E_{pot}(t) = \frac{1}{2} T \int u_x^2 dx$

We define: $E(t) = E_{kin}(t) + E_{pot}(t)$



 $E(H) = E_{kln}(H) + E_{pot}(H)$ Proof; $p_{u_{tt}} = T_{u_{xx}} = \int_{-\infty}^{\infty} \mathcal{U}_t \mathcal{U}_{tt} \, dx + T \int_{-\infty}^{\infty} \mathcal{U}_x \mathcal{U}_{xt} \, dx$ $= T \int_{-\infty}^{\infty} \mathcal{U}_t \mathcal{U}_{xx} \, dx + T \int_{-\infty}^{\infty} \mathcal{U}_x \mathcal{U}_{xt} \, dx$ parts of the first = - T J utx ux dx + [T utux] = + T J ux uxt dx term

The first and last terms cancel out, and the middle term is also 0 since u and its derivatives are assumed to vanish "at" $\pm \omega$. Therefore:

 $E'(t) = 0 \implies E(t) = E(0) = const.$

For the plucked string example from before:

$$E = E(0) = \frac{1}{2} T \int_{-\infty}^{\infty} \phi_{X}^{2} dx = \frac{1}{2} T \left(\frac{1}{a}\right)^{2} 2a = \frac{Tb^{2}}{a}.$$
(slope)² length of interval